Chemical Science

Accepted Manuscript

This article can be cited before page numbers have been issued, to do this please use: S. an, D. Kim, J. Kim and S. K. Kim, *Chem. Sci.*, 2023, DOI: 10.1039/D3SC04342H.



This is an Accepted Manuscript, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this Accepted Manuscript with the edited and formatted Advance Article as soon as it is available.

You can find more information about Accepted Manuscripts in the <u>Information for Authors</u>.

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard <u>Terms & Conditions</u> and the <u>Ethical guidelines</u> still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this Accepted Manuscript or any consequences arising from the use of any information it contains.



rsc.li/chemical-science

View Article Online

View Journal

ARTICLE

Received 00th January 20xx, Accepted 00th January 20xx

DOI: 10.1039/x0xx00000x

ce Accepted Manus

Excited-state chemistry of the nitromethane anion mediated by the dipole-bound states revealed by the photofragment action spectroscopy

Sejun An, a Dabin Kima, Junggil Kima and Sang Kyu Kim*a

We report the first experimental observation of the excited dipole-bound state (DBS) of the cryogenically-cooled nitromethane anion (CH₃NO₂⁻), where the excess electron is loosely attached to the singlet or triplet neutral-core. Photofragment and photodetachment action spectra have been employed for the dynamic exploration of Feshbach resonances located even far above the electron detachment threshold, giving the excitation profiles from the ground anionic state (D₀) to the DBSs which match quite well with the spectral structures of the photoelectron spectra. This indicates that the electron transfer from the nonvalence orbital (of DBS) to the valence orbital (of anion) is mainly responsible for the anionic fragmentation channels, giving the strong evidence for that the DBS plays the dynamic doorway-role in the anionic fragmentation reactions. Photofragment action spectra have also been obtained for the anionic clusters of (CH₃NO₂)₂⁻, (CH₃NO₂)₃⁻, or (CH₃NO₂·H₂O)⁻, giving the relative yields of various fragments as a function of the excitation energy for each cluster. The absorption profiles of the anionic clusters exhibit substantial blue-shifts compared to the bare nitromethane anion as their ground states are much stabilized by solvation. The anionic fragmentation pattern varies among different clusters, giving essential clues for the thorough understanding of the whole anionic dynamics such as the dynamic role of the short-lived nonvalence-bound states of the clusters.

Introduction

This article is licensed under a Creative Commons Attribution 3.0 Unported Licence.

Article. Published on 16 October 2023. Downloaded on 10/18/2023 7:18:35 AM.

Open Access

Entry or exit of the electron in the respective form of electronattachment or electron-detachment is the fundamental step in the anion chemistry or physics. Actually, movement of the electron in and out of the anionic system is ubiquitous and essential in a number of chemical¹ and biological² functional processes of redox reactions, dissociative electron attachment (DEA),³ or formation of the interstellar species.^{4, 5} It is quite common that the excess electron in the stable anion occupies one of the vacant valence orbitals of the otherwise neutral molecule, giving the significant changes of the geometrical structure and chemical reactivity whereas the electron affinity estimated from the photoelectron spectroscopy of the anion represents how strongly the electron is bound to the neutral core. On the other hand, since its first conception by Fermi and Teller,⁶ the nonvalence-bound state (NBS) where the excess electron is only loosely bound to the neutral-core without occupying the valence orbital has long been intensively investigated for many recent decades. NBS, according to the nature of the electron binding potential, may be classified into dipole-bound state (DBS),^{7, 8} quadruple-bound state (QBS),⁹ or correlation-bound state (CBS).10 Although the electrostatic potential such as charge-dipole or charge-quadruple interaction is considered to be mainly responsible for the electron binding

of DBS or QBS, respectively, the quantum-mechanical correlation effect and/or the charge-induced-dipole interaction (due to the large polarizability) could also contribute quite significantly in a cooperative way. Especially, recent dynamic¹¹⁻¹⁵ and spectroscopic^{9, 16-21} studies of NBS in the cryogenically-cold condition²² have triggered the molecular-level investigation of the electron-binding nature of NBS.

While the recently-developed time-resolved electrondetachment dynamic studies^{10-15, 23-25} (e.g. autodetachment) at well-defined Feshbach NBS resonances turn out to be extremely useful for disentangling the electron-binding dynamics, the experimental approach to unravel the dynamic role of the NBS as the doorway into various anionic (half- or full-collisional) reactions seems to be still in infancy. In this regard, a very recent work¹³ on the anionic fragmentation reaction mediated by the DBS of ortho-, meta- or para-iodophenoxide is quite noteworthy. Therein, the I⁻ fragment action spectra reflect the quantum resonance structures of DBS whereas the statespecific fragment yields give the efficiency of the anionic fragmentation channel compared to that of the competitive autodetachment process. This may belong to a new form of DEA as the excess electron loosely attached to the positive end of the dipole (of the DBS prepared by the optical excitation) may be transferred into the specific valence orbital leading to either the prompt rupture of the particular chemical bond or rather slow anionic chemical-bond dissociation taking place on the vibrationally-hot anionic state. It is interesting to note that the anionic fragmentation pathway mediated by DBS (prepared by the optical excitation) could be finely controlled in terms of the energy given to the system as well as associated quantummechanical states in contrast to the more traditional electron-

a. Department of Chemistry, KAIST, Daejeon 34141, Republic of Korea. E-mail: sangkyukim@kaist.ac.kr

⁺ Electronic Supplementary Information (ESI) available: [Details of the theoretical calculations for the valence excited states, vibrational structure analysis, and femtosecond resonant two-photon photoelectron spectra of the nitromethane cluster anions]. See DOI: 10.1039/x0xx00000x

ARTICLE

This article is licensed under a Creative Commons Attribution 3.0 Unported Licence.

Open Access Article. Published on 16 October 2023. Downloaded on 10/18/2023 7:18:35 AM.

collision experiment of DEA. In this regard, the photofragment excitation (PHOFEX) spectroscopy where the specific anionic fragment yield is monitored as a function of the excitation energy could be an extremely useful tool for unravelling the nature of the dynamic doorway states into the anionic fragmentation channels, as successfully demonstrated in the recent report on the reaction dynamics of the ortho-, meta- or para-iodophenoxide anions (vide supra). The PHOFEX spectroscopy has great advantages over the photodetachment spectroscopy (where the photoelectron yield is monitored as a function of the excitation energy) in terms of its little background signal as the directly-ejected photoelectron in the latter persists and even increases with increasing the excitation energy above the electron-affinity threshold whereas the absorption profile of the anion species is supposed to be straightforwardly reflected in the PHOFEX spectroscopy. Furthermore, it would be even ideal for the investigation of the excited-state chemistry over the wide excitation energy region even far above the detachment threshold,²⁶⁻²⁹ if and only if it is associated with any fragmentation channel in the overall anionic reaction. Despite these advantages (compared to other conventional photoelectron and photodetachment spectroscopic methods), the PHOFEX spectroscopy has been surprisingly rarely employed for chemical dynamic study of the anion, mainly because of the extremely low yields of the anionic fragments at the excitation energy above the photodetachment threshold.

Herein, we have investigated the excited-state chemistry of the nitromethane anion (CH₃NO₂-) using the PHOFEX spectroscopy in addition to the photoelectron and photodetachment spectroscopic methods. The nitromethane anion is one of the most-studied radical anion systems³⁰⁻³⁶ whereas the presence of its DBS, as the dipole moment of nitromethane is 3.46 Debye,³⁰ has been demonstrated on several occasions beforehand. Probably due to the very short lifetime, however, Feshbach resonances of the DBS of nitromethane anion have not yet been observed. As a matter of fact, from the (I·CH₃NO₂)⁻ cluster study, the Neumark group^{25, 37} reported that the nitromethane DBS anion survives only briefly ($\tau \sim 400$ fs) before it relaxes into the valence anionic states. The direct excitation spectrum to the DBS from the anionic ground state thus has not been reported to date, hampering the further reaction dynamics study of the current system. In this work, by employing the PHOFEX spectroscopy, we were able to unravel the absorption profiles reaching to the DBSs associated with the ground (S0) and several excited (T1, S1, T2) states of the neutralcore over the wide excitation energy range as well as the statespecific fragmentation patterns for the first time. PHOFEX spectroscopic study for the anionic clusters of $(CH_3NO_2)_2$, $(CH_3NO_2)_3$, or $(CH_3NO_2 \cdot H_2O)$ gives the unprecedented information regarding the energy flow leading to the solvent evaporation and/or chemical-bond dissociation, which is most likely mediated by the electron conveyance among nonvalence and valence orbitals of the anion. The anion PHOFEX spectroscopy turns out to be a truly excellent experimental tool for unravelling the anionic chemical dynamics especially regarding anionic fragmentation reactions resulting from the

Experimental and computational methods

Experimental methods.

To generate nitromethane cluster anions, neon gas was mixed with liquid nitromethane and its clusters, and then expanded through a pulsed Even-Lavie valve coupled to a filament ionizer. Secondary electrons were produced by the electron ionization of the neon molecular beam and were attached to the nitromethane and its clusters. The generated anions passed through a skimmer into a quadrupole deflector that guides the anions orthogonally and subsequently through an RF quadrupole mass filter (Ardara Technologies) for initial mass selection. The mass-selected ions entered a quadrupole ion trap held at 25K, where they were stored for 10-49 ms and cooled via collisions with He buffer gas (99.999%). The internally cooled anions were accelerated into the time-of-flight region, and then the mass selected anions were intersected by nanosecond and femtosecond laser pulses in the velocity-map photoelectron imaging spectrometer. Photoproducts (electrons or fragments) generated by the laser pulses were detected by chevron-type microchannel plates (MCPs) backed by a phosphor screen. The photoelectron images were recorded by a CMOS camera, while the photodetachment and PHOFEX spectra were recorded by a photomultiplier tube (PMT). Photoelectron images were then reconstructed by the BASEX.³⁸ Tunable nanosecond laser pulses were generated from an ND: YAG laser system equipped with multiple harmonic generators (NT342, Ekspla) for the photodetachment and PHOFEX spectra. Femtosecond laser pulses, used for resonant two-photon photoelectron spectroscopy, were generated from the Ti: sapphire regenerative amplifier (Legend Elite-P, Coherent), which was seeded by the femtosecond oscillator (Vitara-T-HP, Coherent).

Computational details

The ground states geometries of anions (D₀) and neutrals (S₀) were optimized to obtain the thermodynamic threshold energies for each dissociation reaction channel using the density functional theory (DFT) with the combination of B3LYP functional and the 6-311++G(3df,3pd) basis set. The vertical excitation energies of $D_0 \rightarrow D_n$ transitions were calculated using the time-dependent density functional theory at the same computation level as above. The first electronic excited-state (D₁) of the nitromethane anion were optimized using the resolution-of-the-identity second-order approximate coupledcluster singles and doubles (RICC2) method with aug-cc-pVTZ basis set. In order to illustrate the C-N bond dissociation reaction of nitromethane anion, potential energy curves of D₀ and D₁ state were obtained by scanning the C-N bond length while the other molecular geometries are kept frozen at the D₀ equilibrium geometry. All DFT or RICC2 calculations were performed in Gaussian 09³⁹ or Turbomole 7.0.2⁴⁰ program package.

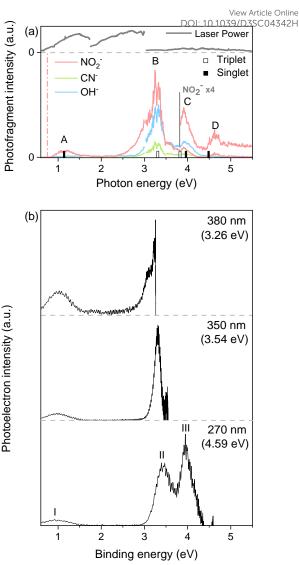
Chemical Science

Chemical Science

Result and Discussion

PHOFEX spectra of the cryogenically-cooled nitromethane anion, taken for the anionic fragments of NO2⁻, OH⁻, and CN⁻ as a function of the excitation energy, are given in Fig. 1. The PHOFEX spectrum of NO2⁻ reflecting its fragmentation yield from the optical excitation of CH₃NO₂⁻ gives the four distinct bands of A, B, C, or D centered on 1.04, 3.26, 3.95, or 4.68 eV, respectively. PHOFEX spectra of OH⁻ and CN⁻ exhibit similarly shaped bands of B, C, and D whereas the A band is observed only for NO_2^- , indicating that the thermodynamic appearance threshold for the OH⁻⁴¹ or CN⁻ fragment should be substantially higher than that for the NO2⁻ fragment. As the PHOFEX signal of the F-fragment (S_F) as a function of the excitation energy (E) is supposed to be proportional to the absorption cross-section of the parent anion (σ_P) times the relative yield of F (P_F), giving the relation of $S_F(E) \propto \sigma_P(E) \cdot P_F(E)$, overall shapes of the individual PHOFEX bands should be governed by the absorption profiles of the parent anion unless the relative yields of the particular anionic fragments vary drastically within the same absorption band.

It should be noted that the electron affinity (EA) of CH₃NO₂is quite low, as it has recently been estimated to be $\sim 0.172 \text{ eV}$,⁴² and thus the photodetachment liberating the electron from the anion upon the optical excitation should be the most efficient process. And yet, as the photodetachment (or photoelectron) yield is the consequence from the optical excitation to the deionization continuum and also because it is cumulative with increasing the excitation energy, the Gaussian-shaped PHOFEX bands (Fig. 1) cannot be attributed to the dynamic behavior of the photodetachment process. The first impression might be then that the A-band (centered at 1.04 eV) could be due to the excitation to the first electronically excited-state (D₁) of CH₃NO₂⁻. Despite that the possible contribution of the first-excited nitromethane anionic state to the A-band may not be completely excluded, there are several evidences against such a scenario. Namely, the guantum-mechanical calculations for the valence excited-states of the nitromethane anion do not reproduce the overall spectral structure of the PHOFEX spectrum (see ESI⁺). The PHOFEX bands could be then due to the optical transition to the dipole-bound states associated with the ground (A) and excited⁴³ (B-D) states of the nitromethane neutral-core. It should be emphasized that the presence of DBS^{30, 32, 36, 42} as well as its ultrafast relaxation^{25, 37} into the anionic state of the nitromethane anion has already been reported in previous other experimental works, suggesting that the optical transition to DBS may end up with the efficient relaxation into the valence states of CH₃NO₂⁻, eventually leading to anionic fragmentation into the CH₃ + NO₂⁻ products, for instance. Namely, as the DBS associated with the ground neutral-core state survives only briefly with the lifetime of ~ 400 fs,^{25, 37} the relatively slow autodetachment process should be kinetically less favored while the ultrafast internal conversion of the DBS into the anionic valence state becomes facilitated. Accordingly, the spectral features of Feshbach resonances of the lowest DBS are not expected to stand out in the photodetachment spectroscopy (Fig. 2), although those are not



cross the range r configuration, rge fluctuations d in the upper e column) and meshold for the tron spectra of 70 nm cited DBSs.

Fig. 1 (a) PHOFEX spectra of cryogenically-cooled nitromethane anion across the range of 0.6 – 5.5 eV. The spectra were acquired in four segments due to laser configuration, with divisions at a photon energy of 1.75, 3.00. and 3.50 eV. Due to large fluctuations in laser power within our spectral window, laser power is represented in the upper panel. The theoretical vertical detachment energies to triplet (white column) and singlet (black column) states are depicted. The theoretical dissociation threshold for the NO₂⁻ fragment is indicated with a vertical dashed line. (b) Photoelectron spectra of nitromethane anion recorded at wavelengths of 380 nm, 350 nm, and 270 nm

anticipated to be identifiable either for the core-excited DBSs. The PHOFEX spectrum where the NO_2^- yield represents the excitation profile to the DBS may then carry the vibronic structures of the meta-stable resonant states. Indeed, in the PHOFEX-A band (also in B-band; *vide infra*), many sharp resonant spectral features could be identified (Fig. 2) though the appropriate mode assignments are nontrivial as those are highly congested due to the much higher appearance energy of NO_2^- (~ 0.75 eV) compared to the EA threshold of 0.172 eV (*vide supra*).

Another strong evidence that the PHOFEX spectrum represents the $D_0 \rightarrow DBS$ excitation is the experimental finding that the overall absorption profile is very similar to that of photoelectron spectrum. It should be noted that the PHOFEX-A

Chemical Science

ARTICLE

This article is licensed under a Creative Commons Attribution 3.0 Unported Licence.

Open Access Article. Published on 16 October 2023. Downloaded on 10/18/2023 7:18:35 AM.

band is apparently truncated at the red-edge of the absorption profile due to the much higher NO2⁻ appearance threshold than the electron detachment threshold (vide supra). Actually, the photoelectron spectra of the cryogenically-cooled nitromethane anion obtained at the excitation energy of 380, 350, or 270 nm (Fig. 1) give three distinct photoelectron bands centered at the binding energy of 0.97, 3.39, and 3.95 eV, which are in excellent agreement with the PHOFEX A, B, and C bands, respectively. In fact, those photoelectron bands of I (0.97), II (3.39) or III (3.95) correspond to the vertical detachment energies (VDEs) estimated for three lowest neutral states of S₀, T₁ and T₂ of CH₃NO₂, respectively, according to the report by Sanov and coworkers,³⁴ although the photoelectron band associated with S1 near the photoelectron band-II (Fig. 1) seems to be nearly absent in the present photoelectron spectrum. It should also be noted that all PHOFEX bands match very well with theoretically-predicted vertical detachment energies (Table I). As the DBS is (intrinsically) nearly identical to the neutral-core in terms of the molecular structure and energetics,19 the fact that the spectral pattern of the PHOFEX spectrum matches with that of the photoelectron spectrum strongly indicates that the former originates from the excitation to the DBS in the wide excitation energy range (0.5 - 5 eV). It should be then followed by the internal conversion to the lowlying valence states of the parent anion eventually leading to NO₂⁻, OH⁻ or CN⁻ fragments. For the more accurate assignment of the PHOFEX-B band, for instance, the more thorough comparison between PHOFEX and photoelectron spectra has been made in terms of the detailed vibrational structure (Fig. 2). The highly-resolved photoelectron band-II from Ref. 34 where the Sanov group had assigned as the triplet state (T1) of nitromethane has been found to be in excellent agreement with the PHOFEX-B band, exhibiting the clear progressions and combinations of ~ 240 and/or ~ 540 cm⁻¹ vibrational modes (see ESI⁺). Particularly, the 540 cm⁻¹ mode assigned as the most Franck-Condon active NO₂ bending matches quite well with the theoretically prediction of 547 cm⁻¹. ³⁴ The PHOFEX-C band associated with the photoelectron-III band may be then attributed to S_1 and/or T_2 of nitromethane, whereas the PHOFEX-D band might reflect S2 (Table I) though it is not unambiguous as the corresponding photoelectron spectrum is not available at the present time.

It should be noted that the determination of the D_0 -DBSs transition oscillator strengths or the fragmentation yields are nontrivial at the present time. Nevertheless, experimental observation of the core-excited DBSs where the excess electron is attached to the excited-states of the neutral-core is quite

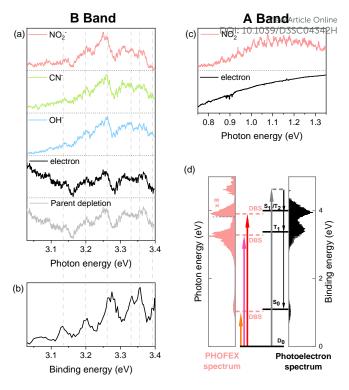


Fig. 2 (a, c) PHOFEX and photodetachment spectra of cryogenically-cooled nitromethane anion in the B band (a) and A band (c) region, respectively, acquired while scanning the photon energy in 4 cm⁻¹ step. The photodepletion spectrum is also recorded for the B band. Signal intensities are normalized for clarity. (b) The partial 355 nm photoelectron spectrum of nitromethane anion in the band II region (3.03-3.40 eV) by Sanov *et al.* (Ref 34). (d) Schematic energy-level diagram showing the anionic ground state, DBSs, and neutral states for nitromethane. The PHOFEX spectrum measured with NO₂⁻ fragments and photoelectron spectrum recorded at 270 nm are plotted. Below the dashed line at 3.26 eV, photoelectron spectrum data is replaced with the 380 nm photoelectron spectrum to accurately describe the lower neutral states (S₀ and T₁).

novel. This also implies that the anionic fragmentation reactions upon the optical excitation, even in the complicated circumstances of the persistence of the detachment continuum, may take place mainly through the DBSs playing as the dynamic doorways. The fragmentation channel responsible for the PHOFEX-A band should then proceed as follows: $CH_3NO_2^-$ (D_0) + $hv \rightarrow DBS \rightarrow CH_3NO_2^-$ ($D_0^*/D_n^*(\sigma_{CN}^*)$) $\rightarrow CH_3 + NO_2^-$. The bond dissociation would occur either statistically on D_0^* (slow) or promptly on the optically-dark $D_n^*(\sigma_{CN}^*)$ state (fast). It has been found that many valence excited-states of the nitromethane anion are very crowded in the narrow excitation energy range, and it seems to be quite challenging to designate any particular excited-state responsible for the specific fragmentation channel.

Photoelectron band	Neutral state	Experimental D₀→neutral (eV)	Theoretical D₀→neutral (eV)	PHOFEX band	Experimental $D_0 \rightarrow DBS (eV)$
I	So	0.97	1.14	А	1.04
Ш	T_1	3.39*	3.31	В	3.26
III	T ₂ /S ₁	3.95/4.12ª	3.82/3.96	С	3.95
	S ₂		4.49	D	4.68

This journal is © The Royal Society of Chemistry 2023

This article is licensed under a Creative Commons Attribution 3.0 Unported Licence.

Open Access Article. Published on 16 October 2023. Downloaded on 10/18/2023 7:18:35 AM.

Chemical Science

And yet, the presence of $D_n^*(\sigma_{CN}^*)$ which is repulsive with respect to the C-N bond elongation coordinate could be confirmed from our potential energy surface calculations (see ESI⁺). The NO₂-fragmentation from D_0^* (slow) would be less likely as the relatively faster vibrational autodetachment⁴⁴ from D_0^* should be kinetically dominant. In this sense, the NO₂⁻ fragment from the PHOFEX-A band is most likely due to the ultrafast bond rupture taking place eventually on the σ_{CN}^* state. All of NO_2^- , OH^- , and CN^- fragments are produced from the PHOFEX-B, C, D bands, and it may result from the vibronic couplings among various electronic states of the anion. Incidentally, it is intriguing to note that the anionic fragmentation reaction originated from the core-excited DBS had not been reported to date. In this regard, the presence of the high-lying DBSs manifested by the PHOFEX-B, -C, or -D band strongly suggests that the DBS plays the important role as the doorway into the nonadiabatic transition even to the anionic excited-states, leading to the anionic fragments. Li et. al. recently reported that the electron of the high kinetic energy range (4 - 8 eV) could be attached to the neutral peptide molecules (leading to the fragmentation) in the situation where the corresponding anionic valence states are hardly accessible.45 Accordingly, their electron-attachment spectrum was attributed to DEA possibly mediated by the core-excited DBSs. Our experimental finding supports such a scenario that the core-excited DBS may play the important dynamic role for the DEA occurring in the high collisional energy region. Besides, dynamics of the anionic fragmentation reactions associated with PHOFEX-B, C, D bands are certainly subject to the further investigation.

For the anionic clusters of (CH₃NO₂·H₂O)⁻, (CH₃NO₂)₂⁻, or $(CH_3NO_2)_3^-$ the PHOFEX spectrum of each species gives the excitation profiles quite similar to those of the nitromethane anion in the 0.55 – 3.0 eV range (Fig 3). Anionic fragments from the $(CH_3NO_2)_2^-$ dimer, for instance, consist of $CH_3NO_2^-$, $(CH_3NO_2) \cdot NO_2^-$, or NO_2^- which result from the solvent (CH_3NO_2) evaporation, C-N bond rupture, or the combination of these, respectively. The PHOFEX excitation profiles of those anionic fragment species are different in terms of their appearance/disappearance energies. The CH₃NO₂- PHOFEX band (due to the evaporation of the solvent) from the dimer anion starts to appear at ~ 0.8 eV whereas its asymmetric-Gaussian-shaped excitation profile diminishes at ~ 1.7 eV. The appearance energy of the CH₃NO₂⁻ fragment coincides with that of the photoelectron, suggesting that the corresponding PHOFEX signal may indeed originate from the excitation to the NBS of the anion dimer. It should be noted that the electron

Table 2 Vertical transition energies for the $D_0{\rightarrow}S_0$ and $D_0{\rightarrow}NBS$ of the nitromethane clusters

Cluster	Theoretical D₀→S₀ (eV)	Experimental ^a D₀→S₀ (eV)	Experimental D₀→NBS (eV)
(CH₃NO₂·H₂O) ⁻	2.0	1.7	1.6
(CH ₃ NO ₂) ₂ -	1.9	1.7	1.6
(CH ₃ NO ₂) ₃ ⁻	2.4	2.2	2.2
^a Reference [35]			



ARTICLE

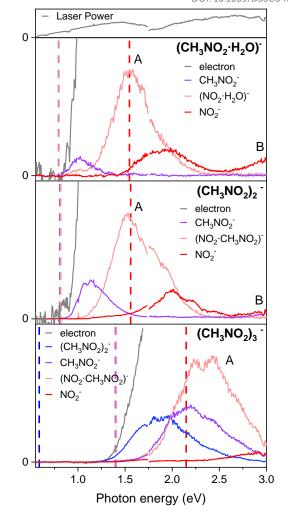


Fig. 3 The PHOFEX and photodetachment spectra of cryogenically-cooled (CH₃ NO₂·H₂O)⁻, (CH₃NO₂)₂⁻, and (CH₃NO₂)₃⁻. Due to the laser configuration, the spectra were acquired in two segments with a division at photon energy of 1.75 eV and the laser power is represented at the top. The theoretical dissociation threshold values for each fragment are indicated with vertical dashed lines of the corresponding colors. Details of the calculations are described in ESI.

affinity of $(CH_3NO_2)_2^-$ is much higher than that of $CH_3NO_2^-$ as the ground anionic state is much stabilized by the neutral nitromethane solvation.³⁵ The blue-shifted (compared to the monomer) $D_0 \rightarrow NBS$ transition of the dimer anion is thus likely to be responsible for the solvent-solute anion dissociation also because the otherwise direct absorption to the vibrationallyhot D₀* of the anion dimer from the ground state is not expected to be significant. The CH₃NO₂- fragmentation from $(CH_3NO_2)_2^{-1}$ should result from the vibrational excitation of van der Waals modes in the charge-dipole interaction potential, most probably via the nonadiabatic transition from NBS to D₀* of the dimer anion. It should be noted that the photoelectron spectra of the cluster anions reported in the previous other work³⁵ are nearly identical to the PHOFEX spectra obtained in this work, strongly suggesting that the PHOFEX spectra should represent the $D_0 \rightarrow NBS$ transitions of the cluster anions (Table 2). It is interesting to note that the vibrational spectral features

ARTICLE

observed in the PHOFEX-A band of the monomer anion (vide supra) could not be observed in the PHOFEX spectra of the cluster anions. This should be due to the spectral congestion caused by the high density of van der Waals modes of the latter. For the (CH₃NO₂)·NO₂⁻ fragment channel, the Gaussian-shaped PHOFEX profile (1.1 - 2.2 eV) should then represent the $D_0 \rightarrow NBS$ transition followed by the C-N bond rupture on the repulsive σ_{CN}^* state, as similarly found in the case of the monomer anion (vide supra). Therefore, the PHOFEX profile of the $(CH_3NO_2) \cdot NO_2^{-}$ fragment may reflect the coupling strength of the NBS to the repulsive σ_{CN}^* state of the dimer anion as a function of the excitation energy. For the combination channel of the solvent evaporation and C-N bond rupture to give NO2from $(CH_3NO_2)_2^{-}$, as expected from the thermodynamic point of view, its PHOFEX yield shows the appearance threshold of ~ 1.5 eV whereas it diminishes at ~ 2.5 eV. The NO_2^- fragment from the anion dimer could be the consequence from the sequential process of the C-N bond rupture followed by the solvent evaporation, although the detailed mechanism is subject to the further investigation. The similar systematic approach could be applied for the anionic cluster of $(CH_3NO_2)_3^-$ or $(CH_3NO_2 \cdot H_2O)^-$, and the more detailed dynamics study on those anion cluster systems will be forthcoming soon.

For the anion cluster of (CH₃NO₂)₂⁻ or (CH₃NO₂)₃⁻, for example, it could be arguable if its DBS could exist as the dipole moment of the nitromethane dimer (as the neutral-core) may be far below the conventional threshold value of 2.5 D^{18, 46} required for the presence of DBS. It should be though emphasized that there have been a significant number of cases where other types of NBS such as CBS have been identified for the anionic cluster species.^{10, 24} Therefore, it is highly plausible for the anionic clusters studied here to be present as NBS. In order to validate such a scenario, we have performed the resonant two-photon photoelectron spectroscopy using the femtosecond (fs) laser pulse. Because of the ultrashort lifetime the NBS species, the photoelectron from the autodetachment of NBS should be hardly detectable (vide supra). Using the fs laser pulses ($\Delta t \approx 100$ fs), we could clearly observe the anisotropic photoelectron band with the zerobinding energy from the two-photon excitation of the $(CH_3NO_2)_2^-$ or $(CH_3NO_2)_3^-$ at the respective photon energy of ~ 0.83 eV or ~ 1.55 eV (see ESI⁺). Here, the first photon excites the ground anion to the NBS whereas its loosely-attached excess electron is detached by the second photon within the same fs laser pulse. This experimental finding strongly indicates that the NBS of the nitromethane dimer or trimer anion actually exists. From the fact that the dipole moments of the cluster species of nitromethane are quite low, the correlation effect is most probably responsible for the electron binding of the NBS species of these cluster anions.

Conclusions

The anion PHOFEX spectroscopy has been employed to give the first experimental observation of the ground and core-excited dipole-bound states of the cryogenically-cooled nitromethane anion (CH_3NO_2) which are associated with the S₀, T₁ or T₂ state

of the neutral-core. The excitation profiles reflected icin the PHOFEX spectra represent the $D_0 \rightarrow DBS$ transitions/responsible for the anionic fragmentation reactions, indicating that the DBS may play the doorway role into the anionic fragmentation reaction leading to NO₂⁻, OH⁻, or CN⁻. The overall spectral features of the PHOFEX bands are in excellent agreement with those of the photoelectron bands, strongly indicating that the transfer of the excess electron in the nonvalence orbital (of DBS) to the valence orbital (of anion) should be largely responsible for the anionic fragmentation channels. Photofragment action spectra for the anionic clusters of (CH₃NO₂)₂⁻, (CH₃NO₂)₃⁻, or (CH₃NO₂·H₂O)⁻ give essential information regarding the complicated fragmentation reactions resulting from the electron transfer from the metastable NBS to the excited-states of the cluster anions.

Data availability

The datasets generated and/or analyzed during the current study are available from the corresponding author on request.

Author Contributions

S. A. and D. K. conducted whole experiments. S. A. wrote the paper. J. K. performed computations. S. K. K. conceived the core idea, supervised the whole projects, and edited manuscript.

Conflicts of interest

There are no conflicts to declare

Acknowledgements

This work was supported by the National Research Foundation of Korea under the Project Numbers of RS-2023-00208926 and 2019R1A6A1A10073887.

Notes and references

- M. O. A. El Ghazaly, A. Svendsen, H. Bluhme, S. B. Nielsen and L. H. Andersen, Chemical Physics Letters, 2005, 405, 278-281.
 The Journal of Chemical Physics, 1996, 104, 7788-7791.
- The Journal of Chemical Physics, 1996, 104, 7788-7791.
 L. G. Christophorou, Electron-Molecule Interactions and their Applications, Academic, 1984.
- 4 G. Felix, T. Marek, V. P. Mikhail and P. M. John, The Astrophysical Journal, 2001, 555, 466.
- 5 T. J. Millar, C. Walsh and T. A. Field, Chemical Reviews, 2017, 117, 1765-1795.
- 6 E. Fermi and E. Teller, Physical Review, 1947, 72, 399-408.
- 7 R. D. Mead, K. R. Lykke, W. C. Lineberger, J. Marks and J. I. Brauman, The Journal of Chemical Physics, 1984, 81, 4883-4892.
- 8 T. D. Märk, Angewandte Chemie International Edition in English, 1990, 29, 437-437.
- 9 G.-Z. Zhu, Y. Liu and L.-S. Wang, Physical Review Letters, 2017, 119, 023002.
- 10 J. N. Bull and J. R. R. Verlet, Sci Adv, 2017, 3, e1603106.
- 11 D. H. Kang, S. An and S. K. Kim, Physical Review Letters, 2020, 125, 093001.

This article is licensed under a Creative Commons Attribution 3.0 Unported Licence.

Article. Published on 16 October 2023. Downloaded on 10/18/2023 7:18:35 AM.

Access

Open

Chemical Science

- 12 D. H. Kang, J. Kim, M. Cheng and S. K. Kim, The Journal of Physical Chemistry Letters, 2021, 12, 1947-1954.
- 13 D. H. Kang, J. Kim, H. J. Eun and S. K. Kim, Journal of the American Chemical Society, 2022, 144, 16077-16085.
- D. H. Kang, J. Kim and S. K. Kim, The Journal of Physical Chemistry Letters, 2021, 12, 6383-6388.
 D. H. Kang, J. Kim and S. K. Kim, Chemical Science, 2022, 12.
- 15 D. H. Kang, J. Kim and S. K. Kim, Chemical Science, 2022, 13, 2714-2720.
- 16 H.-T. Liu, C.-G. Ning, D.-L. Huang, P. D. Dau and L.-S. Wang, Angewandte Chemie International Edition, 2013, 52, 8976-8979.
- 17 H.-T. Liu, C.-G. Ning, D.-L. Huang and L.-S. Wang, Angewandte Chemie International Edition, 2014, 53, 2464-2468.
- 18 C.-H. Qian, G.-Z. Zhu and L.-S. Wang, The Journal of Physical Chemistry Letters, 2019, DOI: 10.1021/acs.jpclett.9b02679, 6472-6477.
- 19 G.-Z. Zhu and L.-S. Wang, Chemical Science, 2019, 10, 9409-9423.
- 20 D.-F. Yuan, Y. Liu, C.-H. Qian, Y.-R. Zhang, B. M. Rubenstein and L.-S. Wang, Physical Review Letters, 2020, 125, 073003.
- 21 D.-F. Yuan, Y. Liu, Y.-R. Zhang and L.-S. Wang, Journal of the American Chemical Society, 2023, 145, 5512-5522.
- 22 X.-B. Wang and L.-S. Wang, Review of Scientific Instruments, 2008, 79, 073108.
- 23 J. N. Bull, C. W. West and J. R. R. Verlet, Chemical Science, 2016, 7, 5352-5361.
- 24 J. P. Rogers, C. S. Anstöter and J. R. R. Verlet, Nature Chemistry, 2018, 10, 341-346
- 25 M. A. Yandell, S. B. King and D. M. Neumark, The Journal of Chemical Physics, 2014, 140, 184317.
- 26 K. O. Uleanya and C. E. H. Dessent, Physical Chemistry Chemical Physics, 2021, 23, 1021-1030.
- 27 K. O. Uleanya, R. Cercola, M. Nikolova, E. Matthews, N. G. K. Wong and C. E. H. Dessent, Journal, 2020, 25.
- 28 N. G. K. Wong, J. A. Berenbeim, M. Hawkridge, E. Matthews and C. E. H. Dessent, Physical Chemistry Chemical Physics, 2019, 21, 14311-14321.
- 29 E. Matthews and C. E. H. Dessent, The Journal of Physical Chemistry Letters, 2018, 9, 6124-6130.
- 30 R. N. Compton, H. S. C. Jr., C. Desfrançois, H. Abdoul-Carime, J. P. Schermann, J. H. Hendricks, S. A. Lyapustina and K. H. Bowen, The Journal of Chemical Physics, 1996, 105, 3472-3478.
- 31 31. T. Sommerfeld, Physical Chemistry Chemical Physics, 2002, 4, 2511-2516.
- 32 C. E. H. Dessent, J. Kim and M. A. Johnson, Faraday Discussions, 2000, 115, 395-406.
- 33 H. Schneider, K. M. Vogelhuber, F. Schinle, J. F. Stanton and J. M. Weber, The Journal of Physical Chemistry A, 2008, 112, 7498-7506.
- 34 D. J. Goebbert, K. Pichugin and A. Sanov, The Journal of Chemical Physics, 2009, 131, 164308.
- 35 C. J. M. Pruitt, R. M. Albury and D. J. Goebbert, Chemical Physics Letters, 2016, 659, 142-147.
- 36 G. Liu, S. M. Ciborowski, J. D. Graham, A. M. Buytendyk and K. H. Bowen, The Journal of Chemical Physics, 2020, 153, 044307.
- 37 A. Kunin, W.-L. Li and D. M. Neumark, Physical Chemistry Chemical Physics, 2016, 18, 33226-33232.
- 38 G. M. Roberts, J. L. Nixon, J. Lecointre, E. Wrede and J. R. R. Verlet, Review of Scientific Instruments, 2009, 80, 053104.
- 39 M. J. Frisch, H. B. Schlegel, G. E. Scuseria, M. A. Robb, J. R. Cheeseman, G. Scalmani, V. Barone, G. A. Petersson, H. Nakatsuji, X. Li, M. Caricato, A. V. Marenich, J. Bloino, B. G. Janesko, R. Gomperts, B. Mennucci, H. P. Hratchian, J. V. Ortiz, A. F. Izmaylov, J. L. Sonnenberg, D. Williams-Young, F. Ding, F. Lipparini, F. Egidi, J. Goings, B. Peng, A. Petrone, T. Henderson, D. Ranasinghe, V. G. Zakrzewski, J. Gao, N. Rega, G. Zheng, W.

- Liang, M. Hada, M. Ehara, K. Toyota, R. Fukuda, J. Hasegawa M. Ishida, T. Nakajima, Y. Honda, O. Kitao, Howard, Source and K. Throssell, J. A. Montgomery, Jr., J. E. Peralta, F. Ogliaro, M. J. Bearpark, J. J. Heyd, E. N. Brothers, K. N. Kudin, V. N. Staroverov, T. A. Keith, R. Kobayashi, J. Normand, K. Raghavachari, A. P. Rendell, J. C. Burant, S. S. Iyengar, J. Tomasi, M. Cossi, J. M. Millam, M. Klene, C. Adamo, R. Cammi, J. W. Ochterski, R. L. Martin, K. Morokuma, O. Farkas, J. B. Foresman, and D. J. Fox, Gaussian 09 (Revision D.01). Gaussian Inc., Wallingford CT, 2009.
- 40 TURBOMOLE V7.0.2, a development of University of Karlsruhe and Forschungszentrum Karlsruhe GmbH, 1989-2007, TURBOMOLE GmbH, since 2007; available from http://www.turbomole.com, 2015.
- 41 D. J. Goebbert, D. Khuseynov and A. Sanov, The Journal of Chemical Physics, 2010, 133, 084311.
- 42 C. L. Adams, H. Schneider, K. M. Ervin and J. M. Weber, The Journal of Chemical Physics, 2009, 130, 074307.
- 43 Y.-R. Zhang, D.-F. Yuan and L.-S. Wang, The Journal of Physical Chemistry Letters, 2022, 13, 2124-2129.
- 44 C. L. Adams, H. Schneider and J. M. Weber, The Journal of Physical Chemistry A, 2010, 114, 4017-4030.
- Z. Li, M. Ryszka, M. M. Dawley, I. Carmichael, K. B. Bravaya and S. Ptasińska, Physical Review Letters, 2019, 122, 073002.
- 46 C. Desfrançois, H. Abdoul-Carime, N. Khelifa and J. P. Schermann, Physical Review Letters, 1994, 73, 2436-2439.